Annealing

- Analogy with physical systems
- Relevant for optimization (not integration)
- Powers of probability distributions emphasize the peaks
- If we are looking for a maximum within a lot of distracting peaks, this can help.

Annealing

- Define a temperature T, and a cooling schedule (black magic part)
- Lower temperatures correspond to emphasized maximal peaks.
 - Hence we exponentiate by (1/T).





Continuous versus discrete variables

- Derivatives of continuous distributions can tell you about the structure of your problem.
 - Opportunities for going much faster
- Naive approach is gradient ascent with added stochastic properties
 - Take a step, then perturb the result.
- Typical approach is to link the probability distribution to a potential energy function
 - Follow the system to find low energy (high probability)
 - Stochastic sampling via random momentum
 - An effective example method is Hybrid Monte Carlo

Hybrid Monte Carlo

- References include Andrieu et al. '01 and Neal'93.
- Link the probability distribution to a potential energy function – Alternate stochastic sampling with "dynamics".
 - The dynamics follow the system to find low energy (high probability)
- HMC is an "auxiliary variable sampler"
 - Important trick
 - To sample p(z) we sample p(z,r) or $p(z,r_1,r_2,...)$
 - Ignore the auxiliary variables when we use the samples.



Hamiltonian Dynamics

$p(\mathbf{z}) = \frac{1}{Z_p} \exp(-E(\mathbf{z}))$

We equate \mathbf{z} with position, so $E(\mathbf{z})$ is the potential energy.

High probability \Leftrightarrow Low energy

$$E(\mathbf{z}) = -\log(Z_p) - \log(p(\mathbf{z}))$$

Hamiltonian Dynamics

Recall that the gradient, ∇ , is the vector of partial derivatives.

Recall from physics that force is the negative gradient of potential energy

From before $E(\mathbf{z}) = -\log(Z_p) - \log(p(\mathbf{z}))$

So force is $-\nabla E(\mathbf{z}) = \nabla (\log(p(\mathbf{z})))$

Hamiltonian Dynamics

$$p(\mathbf{z}) = \frac{1}{Z_p} \exp(-E(\mathbf{z}))$$
 and $\nabla E(\mathbf{z}) = \nabla(-\log(p(\mathbf{z})))$

Let **r** be the momentum vector for the system. Denote the kinetic energy by $K(\mathbf{r})$.

 $K(\mathbf{r}) = \frac{1}{2} \|\mathbf{r}\|^2 = \frac{1}{2} \sum_{i} r_i^2 \qquad (W$

(We assume that mass is one).

Hamiltonian Dynamics

 $H(\mathbf{z},\mathbf{r}) = E(\mathbf{z}) + K(\mathbf{r})$ (conserved)

Our distribution with auxilary variables is

$$p(\mathbf{z},\mathbf{r}) = \frac{1}{Z} \exp(-H(\mathbf{z},\mathbf{r}))$$

Hamiltonian Dynamics

 $H(\mathbf{z},\mathbf{r}) = E(\mathbf{z}) + K(\mathbf{r}) \qquad \text{(conserved)}$

We follow \mathbf{z} according to H with a random \mathbf{r}

This can rapidly transport us towards a local minimun (but not to stay indefinitely) thus avoiding random walk.

To follow *H*, we observe that \mathbf{z} (position) changes proportional to \mathbf{r} (momentum), and \mathbf{r} changes proportion to force $(-\nabla E)$.

 $\Delta \mathbf{r} \propto -\nabla E = \nabla \left(\log \left(p(\mathbf{z}) \right) \right)$ $\Delta \mathbf{z} \propto \mathbf{r}$

Following Dynamics

In HMC we follow the dynamics for L time steps of size τ (tunable parameters).

In the "leap frog" method for each τ .

1. Take 1/2 step in **r**.

2. Take a full step in z.

3. Take 1/2 step in **r**.

Following Dynamics

To take a full step in \mathbf{z} .

 $\mathbf{z}(\tau+1) = \mathbf{z}(\tau) + \varepsilon \cdot (\mathbf{r}(\tau))$

(ε is the step size).

Following Dynamics

For L leap frog steps we have.

- Take 1/2 step in r.
 (L-1) times take a full steps in z, then r.
- 3. Take a full step in **z**.
- 4. Take 1/2 step in **r**

Following Dynamics

To take 1/2 step in **r**.

$$\mathbf{r}\left(\tau+\frac{1}{2}\right) = \mathbf{r}\left(\tau\right) + \frac{1}{2}\varepsilon \cdot \left(-\nabla E(\mathbf{z})\right)$$

Following Dynamics

- After L steps of size t, we are at a new point with some bias of being at a lower potential energy (higher probability) and higher momentum.
- Momentum allows us to jump out of wells.

HMC stochastic step

- Typical instantiations sample the momentum variable
- Two common strategies
 - Sample the **r** independently from a Gaussian
 - Sample **r** from a Gaussian using Gibbs
- Note that in both of these cases the proposals are always accepted.

HMC dynamics step acceptance

- If our integration is perfect (i.e., in the limit as t --> 0) then energy is conserved.
 - Thus the value of distribution $p(\mathbf{z},\mathbf{r})$ is the same after the dynamics.
- If we assume no integration errors, we simply accept this step
- If we want to account for error accumulation, we accept the result according to:

$$\min\left(1, \frac{p(\mathbf{z}^*, \mathbf{r}^*)}{p(\mathbf{z}, \mathbf{r})}\right) = \min\left(1, \exp\left(H(\mathbf{z}, \mathbf{r}) - H\left(\mathbf{z}^*, \mathbf{r}^*\right)\right)\right)$$

Putting it all together (A typical vision lab sampler)

- Discrete variables are sampled using (reversible jump) Metropolis Hastings.
- Continuous variables are sampled using stochastic dynamics (essentially hybrid Monte Carlo).
- Discrete variables typically control topology or components
 - The number of components and their type (block, cylinder)
 - How components are connected (branches from a stem)

A typical vision lab sampler

- Randomly proposing structure is too expensive because of the high rejection rate.
- Solution (part one) is to use data driven sampling
 - Proposals are conditioned on distributions computed before we begin using the data
 - For example, the probability of a corner being present in each point in the image.
- Solution (part two) is to delay acceptance
 - Adjust continuous parameters using stochastic dynamics so that the proposed structure is a good fit to the data.

A typical vision lab sampler

- We thus alternate between
 - (1) data driven proposals for new structure (or to switch or kill existing structure)
 - (2) exploring the continuous parameters of the structure
- Additional gains in optimization through having multiple samplers running in parallel exchange information